

CMB Anisotropies

The fundamental measurement is $T(\theta, \phi)$.

Measure the Temp. anisotropies:

$$\frac{\Delta T(\theta, \phi)}{T} = \frac{T(\theta, \phi) - \langle T \rangle}{\langle T \rangle}$$

Since these measurements are on a sphere, expand these fluctuation maps into spherical harmonics:

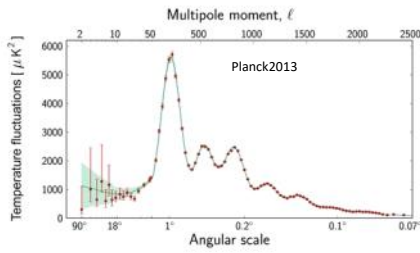
$$\frac{\Delta T(\theta, \phi)}{T} = \sum_{l=1}^{\infty} \sum_{m=-l}^{+l} a_{lm} Y_m^l(\theta, \phi)$$

The coefficients a_{lm} tell us the size of the irreg. on different scales.

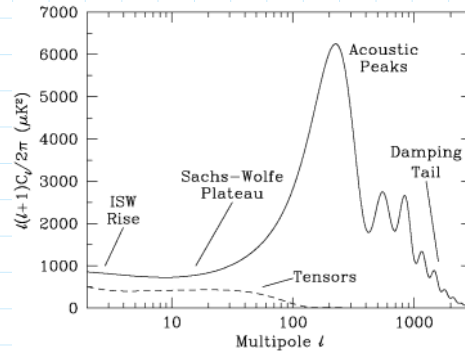
Define the rad'n angular power spectrum

$$C_l = \langle |a_{lm}|^2 \rangle. \text{ The average is over all } m.$$

l can be mapped to angular scale w/ small $l \rightarrow$ large angular scales and large $l \rightarrow$ small angular scales. Rule of thumb, scale $\sim \frac{180^\circ}{l}$



Below $l \sim 200$ is the Sachs-Wolfe plateau due to variations in the DM dist'n at the time of decoupling. The grav-pot variations lead to ΔT variations b/c of grav. redshifts.



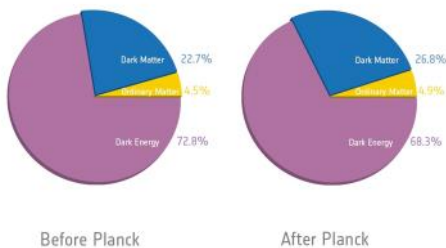
Above $l \sim 200$ the peaks are due to pressure waves in the photon/baryon fluid sloshing in the DM dist'n.

The most important scale is the Hubble time H^{-1} at z_d .

The peak structure comes from oscillations so the 1st peak (at the largest scale) comes from perturbations that had time to ~~etc~~ undergo one oscillation so size scale is cH^{-1} .

Getting the ang. size of that scale depends on d_A at z_d which depends on cosmology, in particular geometry. The location of this first peak strongly suggests $K=0$. Given $\Omega_{m,0} = 0.3$

$$\rightarrow \Omega_{\Lambda} = 0.7$$



The Early Universe

In addition to radiation, neutrinos are another relativistic particle that permeates the Univ.

Assume from now that neutrinos are massless

$$\text{then theoretically } \Omega_{\nu} = 3 \cdot \frac{7}{8} \cdot \left(\frac{4}{11}\right)^{4/3} \Omega_{\text{rad}} \\ = 0.68 \Omega_{\text{rad}}$$

$$\Omega_{\text{rel},0} = \Omega_{\nu,0} + \Omega_{\text{rad},0} = 4.15 \times 10^{-5} h^{-2} \\ \Omega_{\nu,0} = 1.68 \times 10^{-5} h^{-2}$$

$$\text{Since } \rho_{\text{rel}} \propto \frac{1}{a^4} \quad ; \quad \rho_{\text{matt}} \propto \frac{1}{a^3}$$

$$\frac{\Omega_{\text{rel}}}{\Omega_{\text{matt}}} = \frac{4.15 \times 10^{-5}}{\Omega_{\text{m},0} h^2} \cdot \frac{1}{a} \quad \text{where } a_0 = 1$$

So, at decoupling ($a \approx \frac{1}{1000}$)

$$\frac{\Omega_{\text{rel}}}{\Omega_{\text{matt}}} = \frac{0.04}{\Omega_{\text{m},0} h^2} = 0.27 \quad \text{for } \Omega_{\text{m},0} = 0.3 \quad ; \quad h = 0.7$$

So, the Univ. is matter dominated @ decoupling

Radiation/Matter equality occurs at

$$a_{\text{rm}} = \frac{1}{24000 \Omega_{\text{m},0} h^2} = \frac{1}{3528} \quad \text{for } \Omega_{\text{m},0} = 0.3 \quad ; \quad h = 0.7$$

Consider the evol. of T since know $T \propto \frac{1}{a}$

Assume $K=0$, $\Lambda=0$, since small early on

Then $a \propto t^{2/3}$ and $T \propto t^{-2/3}$. Assuming age of Univ. is 12 Gyr (to compensate for ignoring Λ):

$$\frac{T}{2.725 \text{ K}} = \left(\frac{4 \times 10^{17} \text{ s}}{t} \right)^{2/3}$$

$$\text{Then } T_{\text{rm}} = \frac{2.725 \text{ K}}{a_{\text{rm}}} = 66000 \Omega_{\text{m},0} h^2 \text{ K} (= 9614 \text{ K})$$

$$\therefore \text{time of equal. is } t_{\text{rm}} = 10^{11} \Omega_{\text{m},0}^{-3/2} h^{-3} \text{ s}$$

$$\approx 3400 \Omega_{\text{m},0}^{-3/2} h^{-3} \text{ yrs}$$

$$\approx 60000 \text{ yrs}$$

$$(\text{really } 4.7 \times 10^4 \text{ yrs})$$

(compare decoupling @ 350,000 yrs)