

# An Expanding Universe

In 1929, Hubble measured redshifts  $z$  from ~20 galaxies. And then estimated distances to them.

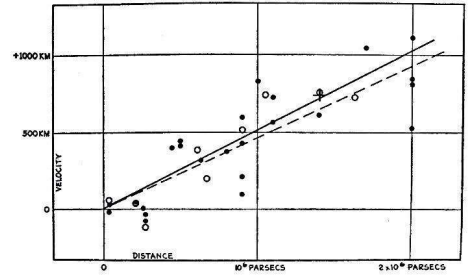


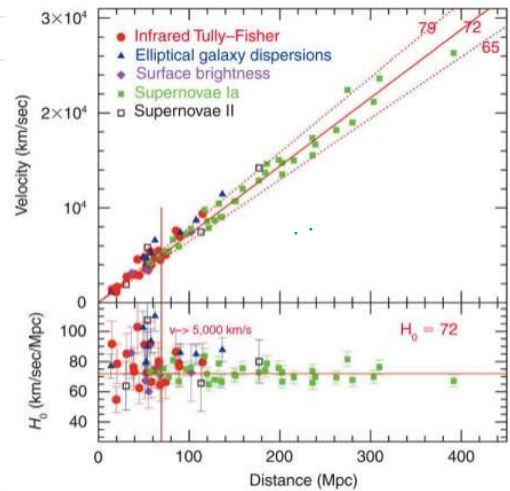
Figure 2.4: Edwin Hubble's original plot of the relation between redshift (vertical axis) and distance (horizontal axis). Note that the vertical axis actually plots  $cz$  rather than  $z$  - and that the units are accidentally written as km rather than km/s. (from Hubble 1929, Proc. Nat. Acad. Sci., 15, 168)

He found  $Cz = H_0 r$   
 where  $H_0$  is a constant (called the Hubble constant)

$$[H_0] = \left[ \frac{1}{s} \right]$$

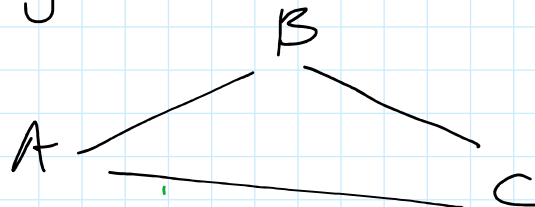
$$\left( \frac{\text{km/s}}{\text{Mpc}} \right)$$

This is exactly what you would expect to see in a Universe undergoing homogeneous and isotropic expansion.



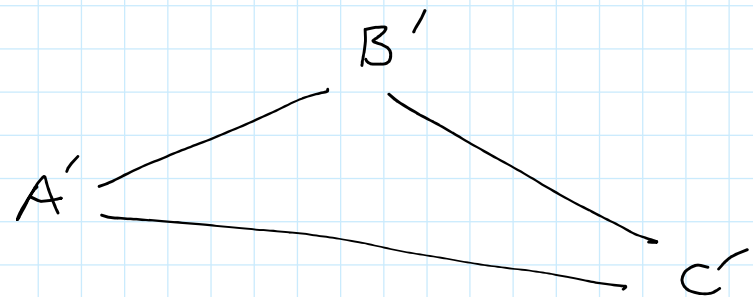
B  
 (Wendy L. Freedman, Observatories of the Carnegie Institution of Washington, and NASA)

Consider 3 galaxies



After homo. and iso. expansion

After homo. and iso. expansion

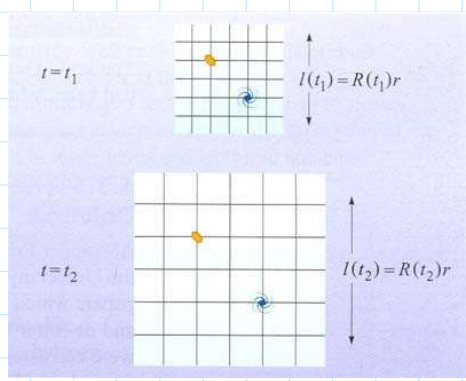


Each side expanded by a constant factor:  
 $a(t) =$  scale factor

Distance from  $A' \rightarrow B'$

$$d_{A'B'} = a(t) l_{AB}$$

$l_{AB} =$  co-moving distance (fixed)



$$\frac{d(d_{A'B'})}{dt} = \frac{da}{dt} l_{AB}$$

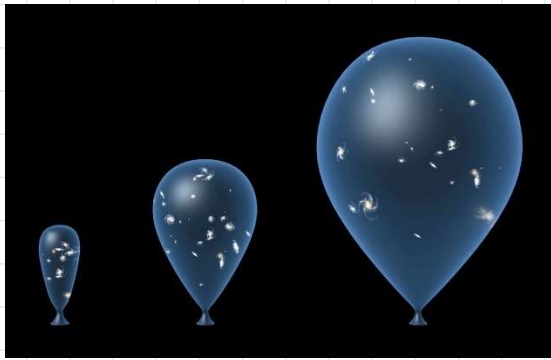
or  $v = \dot{a} l_{AB}$

but  $l_{AB} = \frac{d_{A'B'}}{a(t)} \rightarrow v = \left( \frac{\dot{a}}{a} \right) d$

$v = Hd$

↓ Hubble parameter  
 $H = \frac{\dot{a}}{a}$





$$H = \frac{\dot{a}}{a}$$

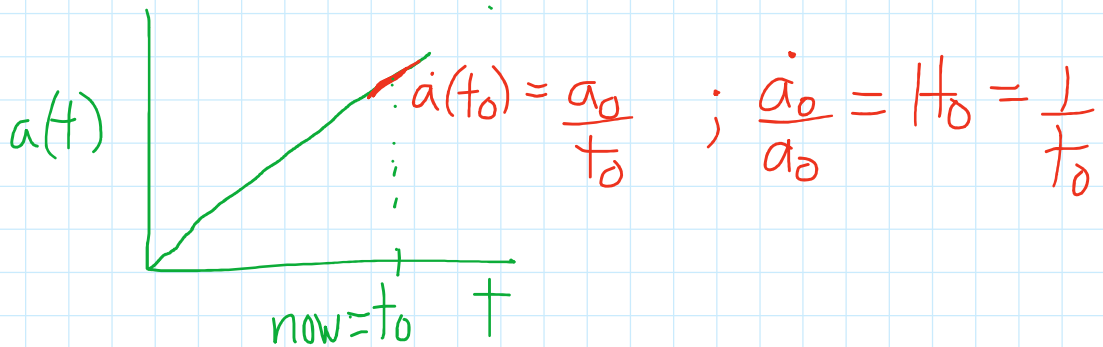
Expanding Univ. means Big Bang!

For small  $v$  ( $v \ll c$ ),  $\frac{\Delta\lambda}{\lambda} = \frac{v}{c}$

or  $v = c \left( \frac{\Delta\lambda}{\lambda} \right) = cz$  where

$$z = \text{redshift} = \frac{\Delta\lambda}{\lambda}$$

If  $\dot{a}(t) = \text{const.}$ , then Hubble parameter now ( $H_0$ ) gives age of Univ.



$$H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$$

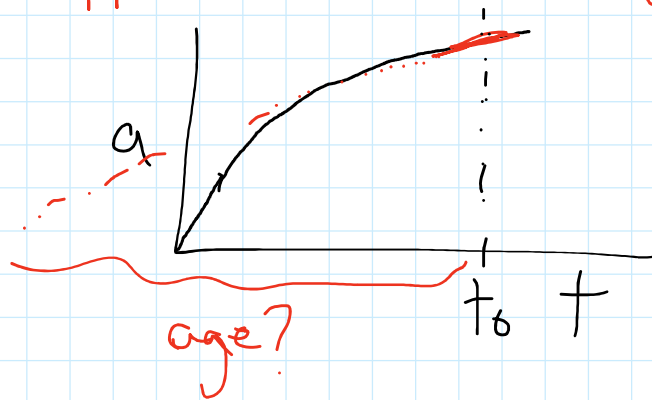
$$\text{so } \frac{1}{H_0} \approx t_0 = 14 \text{ Gyr}$$

('Hubble time')

Similarly, Hubble distance

$$\frac{c}{H_0} \approx 4300 \text{ Mpc}$$

If the Univ. accelerates in any way  
(e.g. due to matter) then  $H_0$  only gives  
an upper-limit to the age:



$$\frac{\dot{a}_0}{a_0} = H_0 < \frac{1}{t_0}$$
$$\hookrightarrow t_0 < \frac{1}{H_0}$$

# An Expanding Universe

In 1929, Hubble measured redshifts,  $z$ , from ~20 galaxies and estimated distances to them. He found

$$cz = H_0 d$$

where  $H_0$  is a constant (called the Hubble constant)

$$[H_0] = \left[\frac{1}{s}\right] \text{ (km/s/Mpc)}$$

This is exactly what you would expect to see in a Universe undergoing homo. and isotropic expansion.

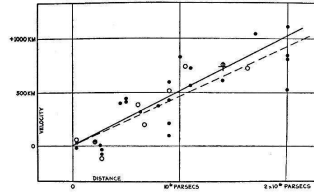
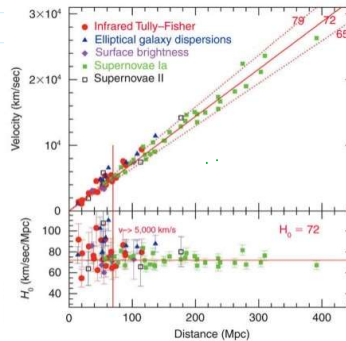
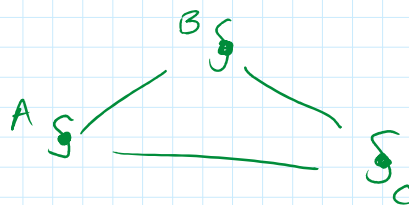


Figure 2.4: Edwin Hubble's original plot of the relation between redshift (vertical axis) and distance (horizontal axis). Note that the vertical axis actually plots  $cz$  rather than  $z$  - and that the units are accidentally written as km rather than km/s. (from Hubble 1929, Proc. Nat. Acad. Sci., 15, 168)

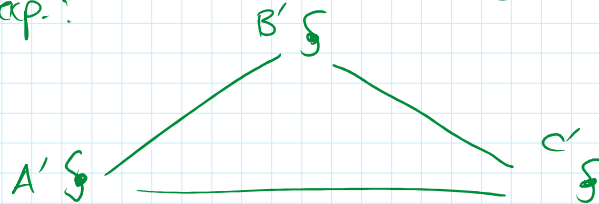


B (Meady L. Freedman, Observatories of the Carnegie Institution of Washington, and NASA)

Consider 3 galaxies:



After iso. & homo. exp.:



Each side has expanded by a const. factor:

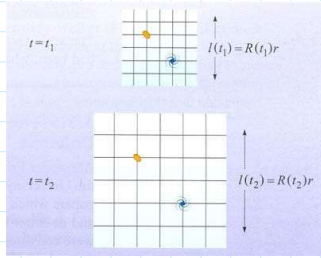
$$a(t) = \text{scale factor}$$

$$\text{Distance from } A' \rightarrow B': d_{A'B'} = a(t) l_{AB}$$

$$l_{AB} = \text{co-moving distance (fixed)}$$

$$d(d_{A'B'}) = d_{A'B'} \dot{a} \quad \text{or} \quad v = \dot{a} l_{AB}$$

$\frac{dL}{dt} = v_{AB}$  or  $v = c \cdot z$



but  $L_{AB} = \frac{dL_{AB}}{dt}$

so  $V = \left(\frac{\dot{a}}{a}\right) d$  → Hubble parameter

$H(t) \equiv \frac{\dot{a}(t)}{a(t)}$

$\therefore V = Hd$

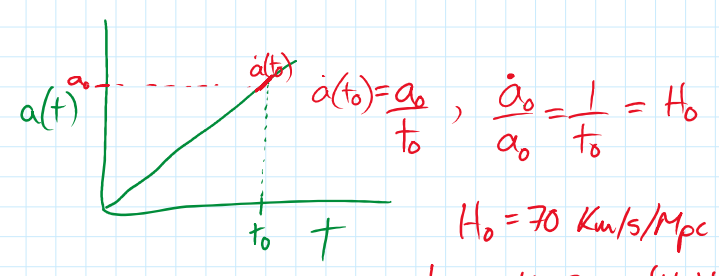
$H_0 = \frac{\dot{a}_0}{a_0}$  now or evaluated at current time



For small  $v$ ,  $\frac{\Delta\lambda}{\lambda} = \frac{v}{c}$  or  $v = cz$

where  $z = \text{redshift} = \frac{\Delta\lambda}{\lambda}$

If  $\dot{a}(t) = \text{constant}$ , then Hubble parameter now ( $H_0$ ) gives age of Univ.

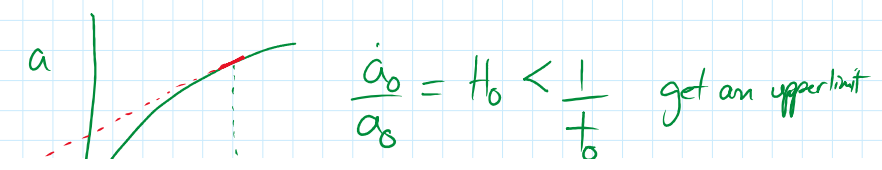


$\dot{a}(t_0) = \frac{a_0}{t_0}$ ,  $\frac{\dot{a}_0}{a_0} = \frac{1}{t_0} = H_0$

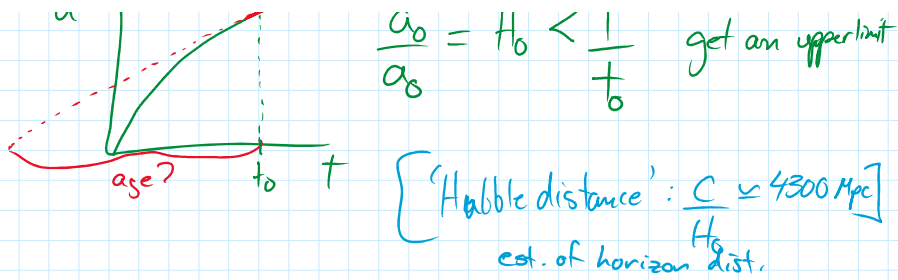
$H_0 = 70 \text{ km/s/Mpc}$

$\frac{1}{H_0} \approx 14 \text{ Gyr}$  'Hubble time'

If the Univ. accelerates in any way (due to matter) then  $H_0$  only gives a limit to the age

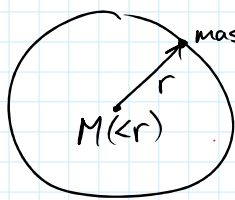


$\frac{\dot{a}_0}{a_0} = H_0 < \frac{1}{t_0}$  get an upper limit



## Newtonian Cosmology

Consider an observer in a uniform expanding medium w/ mass density  $\rho$ . Since the Universe looks the same from anywhere we can consider any point to be its center.



Consider a particle a distance  $r$  away w/ mass  $m$

$$F = m\dot{v} = m\ddot{r}$$

in exp. Univ. (homog; iso)

$$r = a x \quad x: \text{comoving dist}$$

$$\dot{r} = \dot{a} x$$

$$\therefore m\ddot{a}x = -\frac{GmM(<math>r</math>)}{a^2x^2} \quad (\text{no grav. force from outside shell})$$

$$M(<math>r</math>) = \frac{4}{3}\pi r^3 \rho(t) = \frac{4}{3}\pi a^3 x^3 \rho(t)$$

$$\rightarrow \ddot{a}(t) = -\frac{4}{3}\pi G \rho(t) a(t) \quad [\text{Exp. depends only on density.}]$$

$$\text{Cons. of mass: } \rho(t) (a(t)x)^3 = \rho(t_0) (a(t_0)x)^3$$

$$\text{or } \rho(t) = \rho(t_0) \left( \frac{a(t_0)}{a(t)} \right)^3$$

Follow convention  $a(t_0) = 1$

$$\therefore \ddot{a}(t) = -\frac{4}{3}\pi G \rho(t_0) \frac{2\dot{a}}{a^2}$$

$\times$  both sides by  $2\dot{a}(t)$

$$2\dot{a}\ddot{a} = -\frac{4}{3}\pi G \rho(t_0) \frac{2\dot{a}}{a^2}$$

$$\frac{d}{dt}(\dot{a}^2) = +\frac{4}{3}\pi G\rho(t_0)\frac{d}{dt}\left(\frac{1}{a}\right) \cdot 2$$

$$\rightarrow \dot{a}^2 = \frac{8\pi G\rho(t_0)}{3} \frac{1}{a} + \text{const.}$$

$$\therefore \boxed{\dot{a}^2 = \frac{8\pi G\rho a^2}{3} + \text{const.} = \frac{8\pi G\rho a^2}{3} - Kc^2}$$

Friedmann Eqn.

Solving this for  $a(t)$  tells us about the growth, history & future of the Univ.  
Not entirely correct (no  $Gx$ )

Newtonian interp. (mult. both sides by  $\frac{x^2}{2}$ )

$$\frac{(\dot{ax})^2}{2} - \frac{4}{3}\pi(ax)^2\rho G = -kc^2\frac{x^2}{2}$$

$$\frac{v^2}{2} - \frac{4}{3}\pi r^3\rho\frac{G}{r} = -kc^2\frac{x^2}{2}$$

$$\frac{v^2}{2} - \frac{GM}{r} = -kc^2\frac{x^2}{2}$$

$$\left(\frac{KE}{\text{mass}}\right) + \left(\frac{PE}{\text{mass}}\right) = -\frac{kc^2x^2}{2} = \text{const.}$$

Cons. of energy

$K$  & energy of a comoving particle and history of exp. depends on  $K$