

Start again $H^2 = \frac{8\pi G\rho}{3} - \frac{Kc^2}{a^2}$

rewrite this $\frac{a^2 H^2}{c^2} = \frac{8\pi G a^2 \rho}{3c^2} - K$

or $K = \frac{8\pi G a^2 \rho}{3c^2} - \frac{a^2 H^2}{c^2}$

For a given value of H , there is a special density required to make $K=0$. This is the critical density:

$$0 = \frac{8\pi G a^2 \rho_c}{3c^2} - \frac{a^2 H^2}{c^2}$$

$$\therefore \rho_c = \frac{3H^2}{8\pi G}$$

NB: this is a function of time

Currently, $\rho_c(t_0) = 1.88 h^2 \times 10^{-26} \text{ Kg m}^{-3}$
 (compare to H_2O : $\rho = 10^3 \text{ Kg/m}^3$)

Re-write in astro units (check!)

$$\rho_c(t_0) = 2.78 h^2 \times 10^{11} \text{ M}_\odot \text{ Mpc}^{-3}$$

$10^{11-12} \text{ M}_\odot \simeq$ mass of a typical galaxy

$\text{Mpc} \simeq$ typ. galaxy separation

so $\rho \sim \rho_c$ probably...

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ρ_c sets a natural scale for densities in the Univ. \therefore Define the density parameter

$$\Omega(t) = \frac{\rho(t)}{\rho_c(t)}$$

The present value of Ω is Ω_0

$$\text{Now } H^2 = \frac{8\pi G \rho_c \Omega}{3} - \frac{Kc^2}{a^2} = H^2 \Omega - \frac{Kc^2}{a^2}$$

$$\text{or } \boxed{[\Omega - 1] = \frac{Kc^2}{a^2 H^2}}$$

if $\Omega = 1$, $K = 0$ at all times, 'critical univ.'

if $\Omega < 1$, $K < 0$ at all times

if $\Omega > 1$, $K > 0$ at all times

Measuring Ω now (ie, Ω_0) tells us about the fate of the Universe.

If $K \neq 0$, Ω can take any value.

Can have different Ω s, $\Omega_m = \frac{\rho_m}{\rho_c}$, $\Omega_{\text{rad}} = \frac{\rho_{\text{rad}}}{\rho_c}$

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$$\rightarrow \Omega_{dm} = \frac{\rho_{dm}}{\rho_c}, \quad \Omega_r = \frac{\rho_r}{\rho_c}; \quad \Omega = \sum_i \Omega_i$$

often $\Omega_m = \Omega_{bary} + \Omega_{dm}$

The Deceleration Parameter, q_0

Taylor expansion of a around t_0 :

$$a(t) = a(t_0) + \dot{a}(t_0)[t-t_0] + \frac{1}{2}\ddot{a}(t_0)[t-t_0]^2 + \dots$$

\div both sides by $a(t_0)$

$$\frac{a(t)}{a(t_0)} = 1 + H_0(t-t_0) + \frac{1}{2} \frac{\ddot{a}(t_0)}{a(t_0)} (t-t_0)^2 + \dots$$

$$= 1 + H_0(t-t_0) + \frac{1}{2} \frac{\ddot{a}(t_0) H_0^2}{a(t_0) H_0^2} (t-t_0)^2 + \dots$$

Define $q_0 = -\frac{\ddot{a}(t_0)}{a(t_0)} \frac{1}{H_0^2} = -\frac{a(t_0)\ddot{a}(t_0)}{\dot{a}^2(t_0)}$

$$\frac{a(t)}{a(t_0)} = 1 + H_0(t-t_0) - \frac{1}{2} q_0 H_0^2 (t-t_0)^2 + \dots$$

The larger q_0 is, the more rapid the deceleration.

Recall the acc'n eqn: $\frac{\ddot{a}}{a} = -\frac{4\pi G}{3} \left(\rho + \frac{3p}{c^2} \right)$

For matter, $p=0$, $\frac{\ddot{a}}{a} = -\frac{4\pi G \rho}{3}$

So, $\frac{4\pi G \rho_0}{3} \frac{1}{H_0^2}$ but $\rho_{c,0} = \frac{3H_0^2}{8\pi G}$

$$\frac{3}{8\pi G \rho_{c,0}} = \frac{1}{H_0^2}$$

$$\therefore \frac{4\pi G \rho_0}{3} \cdot \frac{3}{8\pi G \rho_{c,0}} = \frac{1}{2} \frac{\rho_0}{\rho_{c,0}} = \frac{1}{2} \Omega_0 \quad (\text{really, } \Omega_{m,0})$$

If we know the prop. of the matter that dominates the Univ, ρ_0 is not independent of H_0 & Ω_0 .

If we don't know, ρ_0 can tell us about acc'n or deceleration indep. of form of energy density.

Age of Univ.

Consider matter dominated $K=0$ Univ. ($\Omega_0=1$)

$$t = \frac{2}{3} \frac{a_m}{c} \left(\frac{a}{a_m} \right)^{3/2}$$

where $a_m = \frac{8\pi G \rho_0}{3c^2}$, so, $t_0 = \frac{2}{3} \frac{a_m}{c} \left(\frac{1}{a_m} \right)^{3/2}$

since $a_0 = 1$ by def'n.

$$\begin{aligned} \therefore t_0 &= \frac{2}{3} \frac{a_m^{-1/2}}{c} \quad \text{but} \quad a_m = \frac{8\pi G \rho_0}{3c^2} \frac{H_0^2}{H_0^2} \\ &= \Omega_0 \frac{H_0^2}{2} = \frac{H_0^2}{2} \end{aligned}$$

$$a = \sqrt{2 \left(\frac{8\pi G\rho_0}{3} \right)^{-1/2} t^2}$$

but from F.E., $H_0^2 = \frac{8\pi G\rho_0}{3}$

$$\therefore a = \sqrt{2H_0 t}$$

at $t = t_0$, $a_0 = 1$,

$$\rightarrow \boxed{t_0 = \frac{1}{2H_0}}$$

