

# Phys 4247 - Assignment #4 Solutions

1. Observable universe constitutes a sphere with proper radius  $d_h = \frac{2c}{H_0}$  for a flat Universe. Since we're dealing with flat geometry, the observable volume is  $V = \frac{4}{3}\pi \left(\frac{2c}{H_0}\right)^3 = \frac{4}{3}\pi \left(\frac{2c}{100 \frac{\text{km}}{\text{Mpc}} h}\right)^3$
- $$\therefore V = 9 \times 10^{11} h^{-3} \text{Mpc}^3$$

Total baryonic mass in the observable Universe:

$$M_b = \Omega_b \rho_{\text{crit}} V = 0.046 \times \left( \frac{2.78 h^{-1} \times 10^{11} M_\odot}{h^{-3} \text{Mpc}^3} \right) \left( 9 \times 10^{11} h^{-3} \text{Mpc}^3 \right)$$
$$= 1.15 \times 10^{22} h^{-1} M_\odot$$

If a typical galaxy has  $10^{11} M_\odot$  of baryonic matter and  $h = 0.7$ , then the # of galaxies in the observable Universe is  $N \approx \frac{1.15 \times 10^{22} (0.7)^{-1} M_\odot}{10^{11} M_\odot / \text{galaxy}} \approx 1.6 \times 10^{11}$  galaxies

Total number of protons. 24% of mass is  ${}^4\text{He}$ , rest is H.

$$\therefore \# \text{ of protons} = 0.24 (1.15 \times 10^{22} h^{-1} M_\odot) \left( \frac{1.989 \times 10^{30} \text{Kg}}{1 M_\odot} \right) \left( \frac{1 {}^4\text{He}}{4 \times 1.67 \times 10^{-27} \text{Kg}} \right) \left( \frac{2 \text{ protons}}{1 {}^4\text{He}} \right)$$
$$+ 0.76 (1.15 \times 10^{22} h^{-1} M_\odot) \left( \frac{1.989 \times 10^{30} \text{Kg}}{1 M_\odot} \right) \left( \frac{1 \text{ H}}{1.67 \times 10^{-27} \text{Kg}} \right) \left( \frac{1 \text{ proton}}{1 \text{ H}} \right)$$
$$= 1.2 \times 10^{79} h^{-1} \text{ protons or } 1.7 \times 10^{79} \text{ protons if } h = 0.7$$

2. a) Density parameter  $\Omega(t) = \frac{\rho(t)}{\rho_c(t)} = \frac{\rho(t)}{3H^2(t)} = \frac{\rho_0 a^{-3} 8\pi G}{3H^2}$

b/c for matter  $\rho(t) = \rho_0 a^{-3}$  (taking  $a_0 = 1$ )

$$\therefore \Omega(t) = \frac{\rho_0 8\pi G}{3H_0^2} \frac{a^{-3} H_0^2}{H^2} = \frac{\rho_0}{\rho_{c,0}} \frac{a^{-3} H_0^2}{H^2} = \Omega_{m,0} \frac{a^{-3} H_0^2}{H^2}$$

but  $(1+z) = \frac{1}{a}$  if  $a_0 = 1$ , so  $\Omega(z) = \frac{\Omega_{m,0} (1+z)^3 H_0^2}{H^2}$

$$\therefore \frac{H^2 \Omega_m(z)}{(1+z)^3} = \text{constant} = \Omega_{m,0} H_0^2$$

b) The Hubble parameter at decoupling as a fraction of its current value is  $\frac{H(z_d)}{H_0} = \frac{\Omega_{m,0}^{1/2} (1+z_d)^{3/2}}{\Omega_m^{1/2}} \approx \Omega_{m,0}^{1/2} (1+z_d)^{3/2}$

where  $z_d = 1100$  is the redshift at decoupling and  $\Omega_m(z_d) \approx 1$

Take  $\Omega_{m,0} = 0.3 \therefore \frac{H(z_d)}{H_0} = \underline{\underline{\approx 2.0 \times 10^4}}$

c) In Assignment #3, Problem #1b, we found that the co-moving distance in a flat Universe with a cosmological constant is

$$r_0 = c H_0^{-1} \int_0^z \frac{dz}{[1 - \Omega_{m,0} + \Omega_{m,0} (1+z)^3]^{1/2}}$$

When  $z \gg 1$ , this becomes  $r_0 = c H_0^{-1} \int_0^z \frac{dz}{\Omega_{m,0}^{1/2} (1+z)^{3/2}} = \frac{c H_0^{-1}}{\sqrt{\Omega_{m,0}}} \left[ \frac{-2}{\sqrt{1+z}} \right]_0^z$

$$= \frac{c H_0^{-1}}{\sqrt{\Omega_{m,0}}} \left( \frac{-2}{\sqrt{1+z}} + 2 \right) \approx \frac{2 c H_0^{-1}}{\sqrt{\Omega_{m,0}}} \text{ if } z \gg 1$$

d) Angle subtended by Hubble length at decoupling is

$$\theta = \frac{cH(z_d)^{-1}}{d_A(z_d)} = \frac{c \Omega_{m,0}^{-1/2} (1+z_d)^{-3/2} H_0^{-1} \Omega_{m,0}^{1/2} (1+z_d)}{2cH_0^{-1}}$$

using the results of parts (b) and (c)

$$\therefore \theta \approx \frac{1}{2} (1+z_d)^{-1/2} = \frac{1}{2} (1101)^{-1/2} = 0.015 \text{ rad} \left( \frac{360 \text{ deg}}{2\pi \text{ rad}} \right) \approx \underline{0.9 \text{ deg}}$$

3.a) The Mattig formula for an open Universe is

$$r_0 = \frac{4c}{H_0} \frac{1}{\Omega_{m,0}^2} \frac{1}{(1+z)} \left[ \frac{\Omega_{m,0} z}{2} + \left( 1 - \frac{\Omega_{m,0}}{2} \right) \left( 1 - (1 + \Omega_{m,0} z)^{1/2} \right) \right]$$

$\therefore$  The angular diameter distance to  $z$  is

$$\begin{aligned} d_{\text{ang}} &= \frac{4c}{H_0} \frac{1}{\Omega_{m,0}^2 (1+z)^2} \left[ \frac{\Omega_{m,0} z}{2} + \left( \frac{2 - \Omega_{m,0}}{2} \right) \left( 1 - (1 + \Omega_{m,0} z)^{1/2} \right) \right] \\ &= 2cH_0^{-1} \frac{(\Omega_{m,0} z + (-1)(\Omega_{m,0} - 2)(-1)(1 + \Omega_{m,0} z)^{1/2} - 1)}{\Omega_{m,0}^2 (1+z)^2} \end{aligned}$$

$$\therefore d_{\text{ang}} = \underline{\underline{2cH_0^{-1} \frac{(\Omega_{m,0} z + (\Omega_{m,0} - 2)((1 + \Omega_{m,0} z)^{1/2} - 1))}{\Omega_{m,0}^2 (1+z)^2}}}$$

b) To use this distance to compute the angular size of the Hubble length at decoupling, split into 2 terms and consider how they behave as  $z$  gets large:

$$d_{\text{ang}} = 2cH_0^{-1} \left[ \frac{\Omega_{m,0} z}{\Omega_{m,0}^2 (1+z)^2} + \frac{(\Omega_{m,0} - 2)(\sqrt{1 + \Omega_{m,0} z} - 1)}{\Omega_{m,0}^2 (1+z)^2} \right]$$



The first term falls as  $\sim 1/z$  as  $z$  gets large. The  $2^{\text{nd}}$  term falls as  $\sim 1/z^{3/2}$  as  $z$  gets large. So at decoupling the  $1^{\text{st}}$  term will dominate and

$$d_{\text{Ang}} \approx \frac{2cH_0^{-1}z}{\Omega_{m,0}(1+z)^2} \approx \frac{2cH_0^{-1}(1+z)}{\Omega_{m,0}(1+z)^2} = \frac{2cH_0^{-1}}{\Omega_{m,0}(1+z)}$$

$$\begin{aligned} \text{So, } \theta &= \frac{cH(z_d)^{-1}}{d_{\text{Ang}}} = \frac{c\Omega_{m,0}^{-1/2}(1+z_d)^{-3/2}H_0^{-1}\Omega_{m,0}(1+z)}{2cH_0^{-1}} \\ &= \frac{1}{2}(1+z_d)^{-1/2}\Omega_{m,0}^{1/2} \approx (1 \text{ deg})\Omega_{m,0}^{1/2} \text{ using the result of 2(d)} \end{aligned}$$

If  $\Omega_{m,0} = 1$ ,  $l \approx 220$ , then this implies

$$\theta \approx \frac{220^\circ}{l}$$

So if  $\Omega_{m,0} = 0.3$ ,  $\theta \approx 0.55 \text{ deg} \rightarrow l \approx 400$

A peak at  $l \approx 400$  is completely incompatible w/ the data.

4. Energy emitted by starlight over 10 Gyr:

$$\begin{aligned} \text{Energy} &= Lt = (2.3 \times 10^{10} L_\odot) \left( \frac{3.85 \times 10^{26} \text{ W}}{1 L_\odot} \right) \left( \frac{3.16 \times 10^7 \text{ s}}{1 \text{ yr}} \right) 10 \times 10^9 \text{ yr} \\ &= \underline{2.8 \times 10^{54} \text{ J}} \end{aligned}$$

If fusion of  ${}^4\text{He}$  produces 28.4 MeV and all the energy in the starlight is generated by fusion then

$$\# \text{ of } {}^4\text{He atoms created} = (2.8 \times 10^{54} \text{ J}) \left( \frac{1 \text{ eV}}{1.602 \times 10^{-19} \text{ J}} \right) \left( \frac{1 {}^4\text{He created}}{28.4 \times 10^6 \text{ eV}} \right) = \underline{6.2 \times 10^{65}}$$

$$\text{Mass in } ^4\text{He created} = 6.2 \times 10^{65} \text{ } ^4\text{He atoms} \left( \frac{4 \times 1.67 \times 10^{-27} \text{ Kg}}{1 \text{ } ^4\text{He}} \right) = 4.1 \times 10^{29} \text{ Kg} = 2.1 \times 10^9 M_{\odot}$$

Original mass in  $^4\text{He}$  using the primordial mass fraction

$$= Y_4 \times 10^{10} M_{\odot} = (0.24)(10^{10} M_{\odot}) = 2.4 \times 10^{10} M_{\odot}$$

New mass in  $^4\text{He} = 2.4 \times 10^{10} M_{\odot} + 2.1 \times 10^9 M_{\odot} = 2.6 \times 10^{10} M_{\odot}$

New  $Y_4 = \frac{2.6 \times 10^{10} M_{\odot}}{10^{10} M_{\odot}} = 0.26$ , so the  $^4\text{He}$  fraction has

increased by 0.02 over its primordial value.

5. The cross-section is  $\sigma_w \sim 10^{-47} \left( \frac{5 \times 10^{-4} \text{ eV}}{10^6 \text{ eV}} \right)^2 \text{ m}^2 = 2.5 \times 10^{-66} \text{ m}^2$

The mean free path of this low energy neutrino through

$^{56}\text{Fe}$  is  $l = \frac{1}{n\sigma_w}$  where  $n$  is the number density of

particles in the lead.

# density of atoms in  $^{56}\text{Fe} = \frac{\rho_{^{56}\text{Fe}}}{M_{\text{nucleus}}} = \frac{7900 \text{ Kg m}^{-3}}{55.8 \times (1.66 \times 10^{-27} \text{ Kg})}$

$$= 8.5 \times 10^{28} \text{ m}^{-3}$$

# density of particles in  $^{56}\text{Fe} = (26 \text{ protons}) 8.5 \times 10^{28} \text{ m}^{-3} + (30 \text{ neutrons}) 8.5 \times 10^{28} \text{ m}^{-3}$

$$+ (26 \text{ electrons}) 8.5 \times 10^{28} \text{ m}^{-3}$$

$$= 7 \times 10^{30} \text{ m}^{-3}$$

$$\therefore l = \frac{1}{(7 \times 10^{30} \text{ m}^{-3})(2.5 \times 10^{-66} \text{ m}^2)} = 5.7 \times 10^{34} \text{ m} = 1.8 \times 10^{12} \text{ Mpc} \gg \frac{c}{H_0}$$