

# Phys 4247 - Assignment #5 Solutions

1. a) Solving for time,  $\left(\frac{1 \text{ sec}}{t}\right) = \left(\frac{T}{2 \times 10^{10} \text{ K}}\right)^2$

$$\therefore t = 1 \text{ sec} \left(\frac{T}{2 \times 10^{10} \text{ K}}\right)^{-2}$$

if  $T = 3 \times 10^{25} \text{ K}$ ,  $t = 4.4 \times 10^{-31} \text{ s}$

b) If  $a(t) \propto t$ , and  $T \propto \frac{1}{a}$  then  $T \propto \frac{1}{t}$

Using result from part (a):

$$T = (3 \times 10^{25} \text{ K}) \left(\frac{4.4 \times 10^{-31} \text{ s}}{t}\right)$$

When  $T$  falls to  $3 \text{ K}$ ,  $t = \frac{(3 \times 10^{25} \text{ K}) \times (4.4 \times 10^{-31} \text{ s})}{(3 \text{ K})}$   
 $= 4.4 \times 10^{-6} \text{ s}$

2. a) Magnetic monopoles behave as non-relativistic matter:

$$\rho_{\text{mon}} = \rho_{\text{mon, GUT}} \left(\frac{a_{\text{GUT}}}{a}\right)^3 = 10^{-10} \rho_{\text{crit, GUT}} \left(\frac{a_{\text{GUT}}}{a}\right)^3$$

Assume Universe is radiation dominated and is at the critical density:

$$\rho_{\text{rad}} = \rho_{\text{rad, GUT}} \left(\frac{a_{\text{GUT}}}{a}\right)^4 = \rho_{\text{crit, GUT}} \left(\frac{a_{\text{GUT}}}{a}\right)^4$$

When these densities are equal:

$$\rho_{\text{crit, GUT}} \left(\frac{a_{\text{GUT}}}{a}\right)^4 = 10^{-10} \rho_{\text{crit, GUT}} \left(\frac{a_{\text{GUT}}}{a}\right)^3$$

$$\text{or } \left(\frac{a_{\text{GUT}}}{a}\right)^4 = 10^{-10} \left(\frac{a_{\text{GUT}}}{a}\right)^3$$

$$\left(\frac{a_{\text{GUT}}}{a}\right)^3 = 10^{-10} \rightarrow \frac{a}{a_{\text{GUT}}} = 10^{10} \rightarrow \left(\frac{t}{t_{\text{GUT}}}\right)^{1/2} = 10^{10}$$

since  $a \propto t^{1/2}$  in the rad'n dominated regime.

From the  $t$ - $T$  relation given in #1,  $\left(\frac{T_{\text{GUT}}}{2 \times 10^{10} \text{ K}}\right) \left(\frac{2 \times 10^{10} \text{ K}}{T}\right) = 10^{10}$

$$\rightarrow T = \frac{T_{\text{GUT}}}{10^{10}} = \frac{3 \times 10^{28} \text{ K}}{10^{10}} = \underline{3 \times 10^{18} \text{ K}}$$

b) Present day:  $\frac{\Omega_{\text{mon},0}}{\Omega_{\text{rad},0}} = \frac{\rho_{\text{mon},0}}{\rho_{\text{crit},0}} \frac{\rho_{\text{rad},0}}{\rho_{\text{crit},0}} = \frac{\rho_{\text{mon},0}}{\rho_{\text{rad},0}} = \frac{\rho_{\text{mon,GUT}} \left(\frac{a_{\text{GUT}}}{a_0}\right)^3}{\rho_{\text{rad,GUT}} \left(\frac{a_{\text{GUT}}}{a_0}\right)^4}$

$$= 10^{-10} \frac{\rho_{\text{crit,GUT}} \left(\frac{a_{\text{GUT}}}{a_0}\right)^3}{\rho_{\text{crit,GUT}} \left(\frac{a_{\text{GUT}}}{a_0}\right)^4} = 10^{-10} \frac{a_0}{a_{\text{GUT}}} = 10^{-10} \left(\frac{t_0}{t_{\text{GUT}}}\right)^{1/2}$$

$$= 10^{-10} \left(\frac{T_{\text{GUT}}}{T_0}\right) = 10^{-10} \left(\frac{3 \times 10^{28} \text{ K}}{3 \text{ K}}\right) = \underline{10^{18} !!}$$

Clearly not compatible w/ observations.

3. For the present density of monopoles to equal that of radiation the density of monopoles has to be reduced by  $\sim 10^{18}$  during inflation (according to the situation of Q#2)

Since  $\rho_{\text{mon}} \propto a^{-3}$  that means  $a$  has to increase by  $\underline{\sim 10^6}$  during inflation.

4. In the development of the ODE for  $\delta$ , there is no dependence on the overall composition or geometry of the Univ. so the basic equation is still  $\ddot{\delta} + 2H\dot{\delta} - \frac{3}{2}\Omega_m H^2 \delta = 0$

Now, consider a negatively curved expanding Univ. w/  $\Omega_m \ll 1$ . In this, as the Univ. is almost 'empty',  $H = \frac{1}{t}$  (Milne Univ.)

$$\therefore \ddot{\delta} + \frac{2\dot{\delta}}{t} = 0$$

Try a power-law solution:  $\delta = Dt^n$

$$n(n-1)Dt^{n-2} + \frac{2Dn}{t}t^{n-1} = 0$$

$$n(n-1) + 2n = 0$$

$$n^2 - n + 2n = n^2 + n = 0$$

$$n(n+1) = 0 \rightarrow n = 0, -1$$

$$\therefore \underline{\delta(t) = D_1 + D_2 t^{-1}}$$

In this Univ, density fluctuations are unable to grow.

5. Need to find radiation and baryonic densities at decoupling  
 $\Omega_{\text{rad},0} = 5 \times 10^{-5}$ ,  $\Omega_{\text{baryon},0} = 4 \times 10^{-2}$

$$\therefore \frac{\Omega_{\text{rad},0}}{\Omega_{\text{baryon},0}} = \frac{5 \times 10^{-5}}{4 \times 10^{-2}} \frac{1}{a} = \frac{1.25 \times 10^{-3}}{a}$$

At decoupling,  $z = 1100$ , so  $a = \frac{1}{1101}$

$$\therefore \frac{\Omega_{\text{rad},0}}{\Omega_{\text{baryon},0}} \approx 1.4 \rightarrow \frac{\rho_{\text{rad},0}}{\rho_{\text{baryon},0}} \approx 1.4 \rightarrow \frac{E_{\text{rad},0}}{E_{\text{baryon},0}} \approx 1.4$$

$$\therefore \epsilon_{\text{baryon}} \approx \frac{\epsilon_{\text{rad}'n}}{1.4} \text{ at decoupling}$$

$$\therefore \epsilon = \epsilon_{\text{rad}'n} + \epsilon_{\text{baryon}} = \epsilon_{\text{rad}'n} + 0.73 \epsilon_{\text{rad}'n} = 1.73 \epsilon_{\text{rad}'n}$$

$$\frac{dP}{d\epsilon} = \frac{dP}{d\epsilon_{\text{rad}'n}} \frac{d\epsilon_{\text{rad}'n}}{d\epsilon} = \frac{1}{1.73} \frac{dP}{d\epsilon_{\text{rad}'n}} = \frac{1}{1.73} \left( \frac{1}{3} \right)$$

$$\therefore \underline{C_s = c \sqrt{\frac{dP}{d\epsilon}} = \frac{1}{\sqrt{1.73}} \frac{c}{\sqrt{3}} = 0.76 \left( \frac{c}{\sqrt{3}} \right)}$$

Since the Jean's Length  $\propto C_s$ , the contribution from the baryons to the Jean's Length at decoupling only decreases it by  $\sim 24\%$  compared to the estimate using just radiation.