

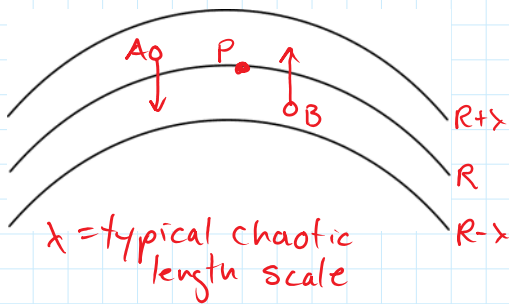
Consider the radius at which the thermal energy per unit mass \approx grav. p.e. per unit mass. At this point, gas will start to slowly accrete to the center

$$\frac{1}{2} C_s^2(\infty) \approx \frac{GM}{r}$$

$$\text{then } r_{\text{acc}} \approx \frac{2GM}{C_s^2(\infty)}$$

Analytical Thin Accretion Disk Theory

Most gas possesses some ang. mom. and will not plunge right into a BH or a NS. Collisions will cancel out the vertical motions but ang. mom. will cause the gas to settle into a disk w/ angular speeds approx. Keplerian: $\Omega_K(R) = \left(\frac{GM}{R^3}\right)^{1/2}$



But because gas particles will jiggle about, viscous stresses are generated which transport ang. mom. orthogonal to the gas motion.

Gas is flowing in ϕ direction
 Gas elements A & B are constantly exchanged across the surface w/ speed $u\vec{v}$ and travel a distance $\sim \lambda$ before interacting

No net transfer of matter, so the mass flux in both directions at P is $\text{Mass flux} \approx \rho\vec{v}$ $\rho = \text{gas density}$

$$\text{Mass rate} \approx \rho \tilde{v} A = \rho \tilde{v} (Hl)$$

where H is the scale height + l is the arc-length along surface

$$\text{So } \frac{\text{Mass rate}}{\text{arc length}} \approx \rho \tilde{v} H$$

But the two masses carry diff. ang. momenta \therefore there will be a viscous torque exerted on the outer stream by the inner stream

Imagine an observer at P rotating w/ $\Omega(R)$, the fluid at $(R - \frac{\Delta}{2})$ will appear to move w/ velocity

$$(R - \frac{\Delta}{2})\Omega(R - \frac{\Delta}{2}) - \Omega(R)R$$

\therefore The average rate of ang. mom./arc length through R in the outward direction is

$$\underbrace{H\rho\tilde{v}}_{\substack{\text{mass rate} \\ \text{arc length}}} \underbrace{(R - \frac{\Delta}{2})}_{(r)} \underbrace{[(R - \frac{\Delta}{2})\Omega(R - \frac{\Delta}{2}) - \Omega(R)R]}_{(v)}$$

inward direction:

$$H\rho\tilde{v} (R + \frac{\Delta}{2}) [(R + \frac{\Delta}{2})\Omega(R + \frac{\Delta}{2}) - \Omega(R)R]$$